

Hadronic Spectral Functions in Lattice QCD at Finite Temperature

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1. Physics of hot/dense QCD
2. Hadrons in hot QCD environment
3. Maximum Entropy Method (MEM)
4. Lattice QCD results
 - mesons on anisotropic lattice at $T \neq 0$
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CP-PACS !

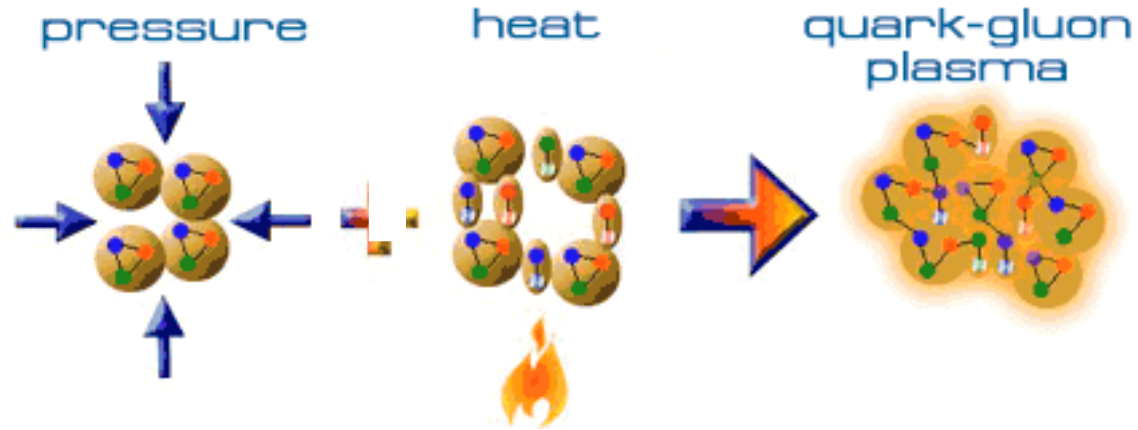
MELQCD Collaboration



M. Asakawa (Kyoto Univ.)
T. Hatsuda (Univ. of Tokyo)
K. Sasaki (Univ. of Tokyo)
S. Sasaki (Univ. of Tokyo)

- Asakawa, Nakahara + T.H., Phys. Rev. D60, 091503 ('99) SR2201
Prog. Part. Nucl. Phys. 46, 459 ('01) (JAERI)
- Asakawa, Nakahara + T.H., hep-lat/0208059 ('02)
- Sasaki, Sasaki, Asakawa +T.H, hep-lat/0209059 ('02) CP-PACS
SR8000(KEK)

高温・高密度における新物質相



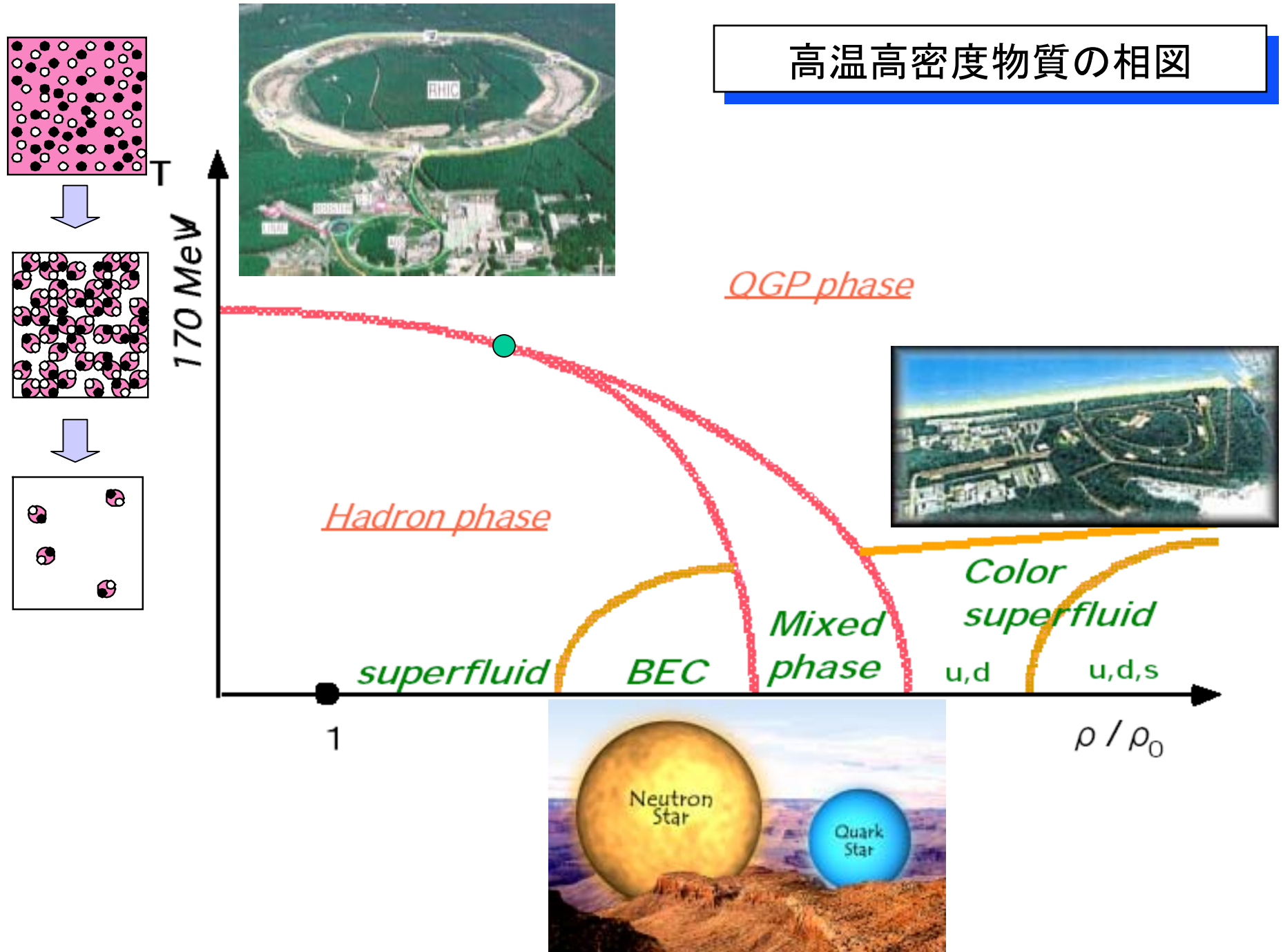
臨界密度: 10^{12} kg/cm^3

臨界温度: 10^{12} K

どこで実現？

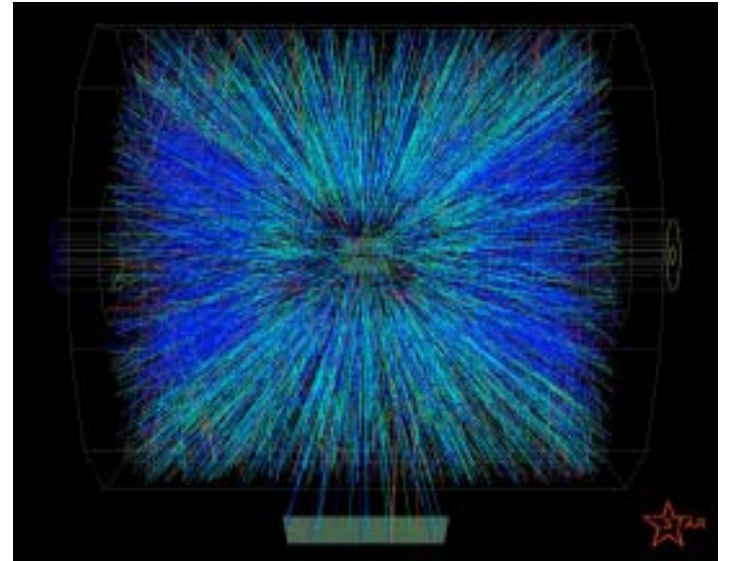
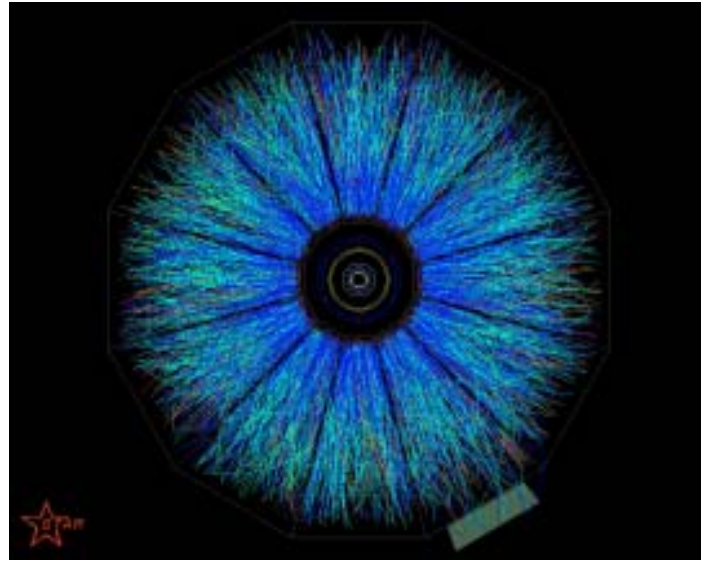
- 初期宇宙 ($t < 10^{-4} \text{ sec}$): 高温・低バリオン密度
- 中性子星中心部: 低温・高バリオン密度
- 相対論的重イオン衝突: 高温 and/or 高バリオン密度

高温高密度物質の相図

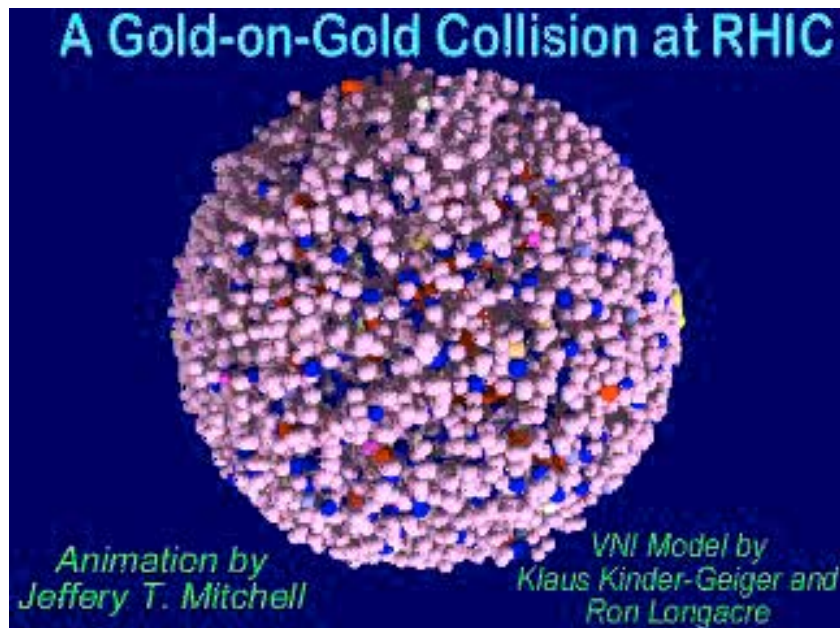


実験室におけるQGP探索

実験 (RHIC)



数値シミュレーション (パートンカスケード)



量子色力学

$$L = -\frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu} + \bar{q} \gamma^\mu (i \partial_\mu - g t^a A_\mu^a) q - m \bar{q} q$$

南部陽一郎 (1966)

Discretization

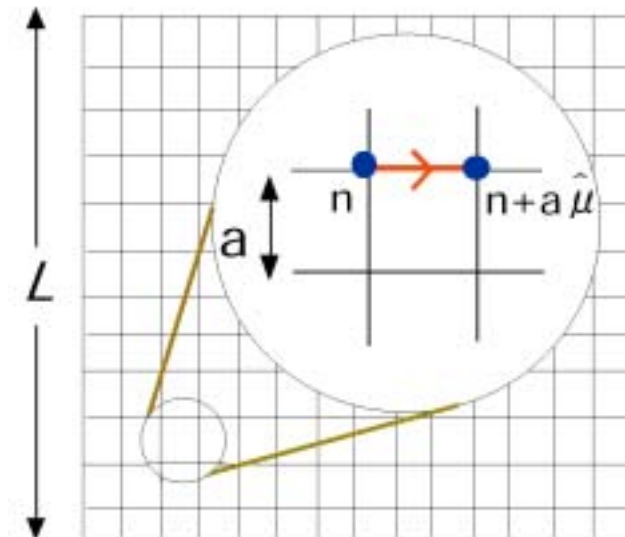
$$A_\mu(x) \rightarrow U_\mu(n) = \exp(i a A_\mu)$$

$$q(x) \rightarrow q(n)$$

$$T=0 : L^3 \times L = (N_s a)^4$$

$$T \neq 0 : L^3 \times 1/T = (N_s a)^3 \times (N_t a)$$

4-d Euclidean lattice



K. Wilson (1973)

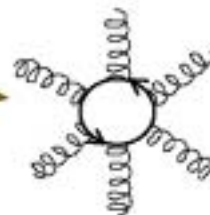
Monte Carlo integration

$$Z = \int [dU] [dq d\bar{q}] e^{-S_{\text{Dirac}}(q, \bar{q}, U) - S_{\text{YM}}(U)}$$

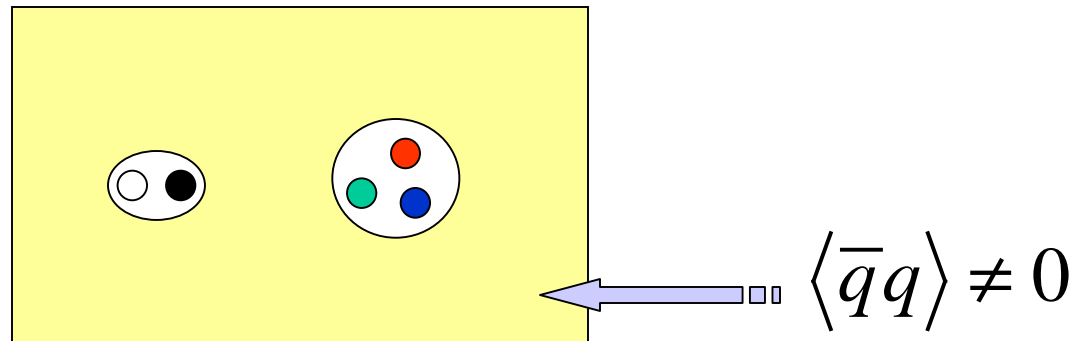
$$= \int [dU] \det Q(U) e^{-S_{\text{YM}}(U)}$$

● Quenched QCD : $\det Q \rightarrow 1$

● Full QCD : $\det Q \neq 1$



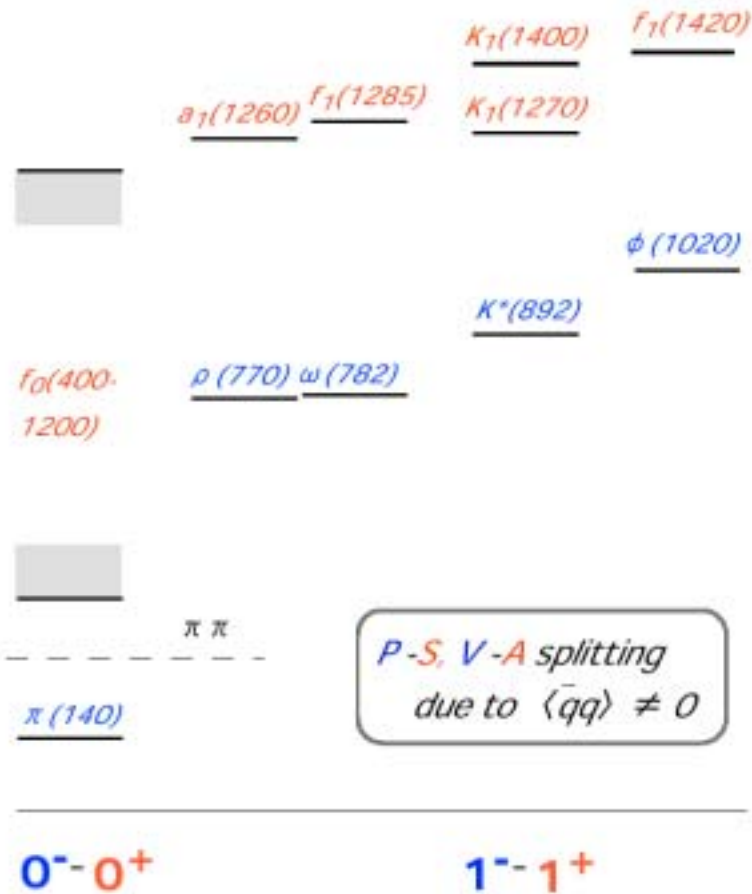
Hadrons in Hot Environment



- 軽いクォーク系とカイラル対称性: *Hatsuda & Kunihiro, PRL ('85)*
- 重いクォーク系と閉じ込め: *Matsui & Satz, PLB ('86)*
Miyamura et al., PRL ('86)

Chiral partners in the vacuum

T.H., soft-dilepton workshop at LBNL ('97)



Tracer of the chiral structure of matter

$$\langle \bar{q} \diamond q(x) \bar{q} \diamond q(y) \rangle$$

- $\langle \bar{q}q \rangle \rightarrow 0$

⇔ Chiral degeneracy

$$\langle S(x)S(y) \rangle \sim \langle P(x)P(y) \rangle$$

$$\langle A(x)A(y) \rangle \sim \langle V(x)V(y) \rangle$$

- change of $\langle \bar{q}q \rangle$

⇔ Individual spectrum

$$\langle S(x)S(y) \rangle, \langle P(x)P(y) \rangle$$

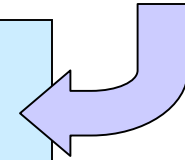
$$\langle A(x)A(y) \rangle, \langle V(x)V(y) \rangle$$

First principle QCD calculation ?

QCD Spectral Functions

Lattice data

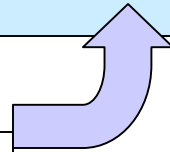
$$D(\tau, \vec{p}) = \int \langle J^+(\tau, \vec{x}) J(0, 0) \rangle e^{i\vec{p}\vec{x}} d^3x$$
$$= \int K(\tau, \omega) A(\omega, \vec{p}) d\omega$$



“Laplace” kernel

$$K(\tau, \omega)$$

$$= e^{-\omega\tau} / (1 \mp e^{-\omega/T})$$



Spectral Function

All information on hadronic correlations
at $T=0$ and $T \neq 0$

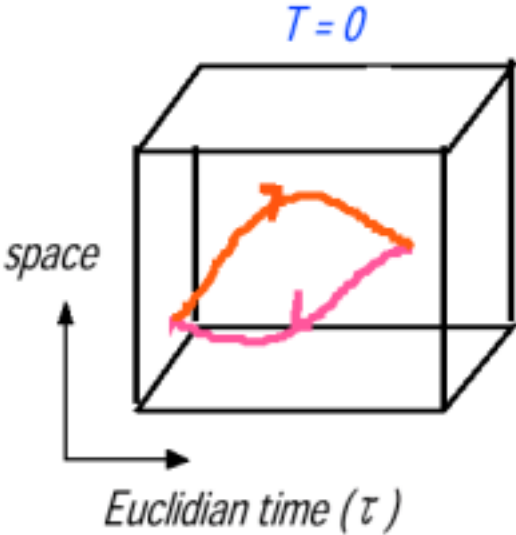
pQCD says ...

$$* A(\omega \geq 0) = \mp A(-\omega) \geq 0$$

$$* A(\omega \gg 1 \text{ GeV}) \rightarrow \omega^{2 \dim[O] - 4} \left(1 + c \frac{\alpha_s}{\pi} \right)$$

How to extract $A(\omega)$ from lattice QCD data?

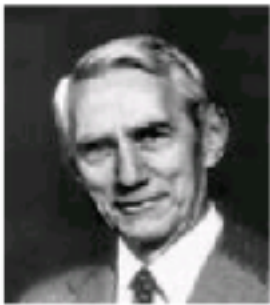
$$\begin{aligned}
 D(\tau) &= \int \langle \bar{q}\Gamma q(\tau, \vec{x}) \bar{q}\Gamma q(0) \rangle d^3x \\
 &\sim e^{-m\tau} \quad (\tau \rightarrow \infty) \\
 &= \int \frac{e^{-\omega\tau}}{1 \mp e^{-\omega/T}} A(\omega) d\omega \quad (\text{any } \tau)
 \end{aligned}$$



Ill-posed problem



T. Bayes 1702-1761



C.E. Shannon, 1916-2001



Maximal Entropy Method : $P[A | D]$

- *No parametrization of $A(\omega)$*
- *unique solution for $D(\tau) \rightarrow A(\omega)$*
- *error estimate on $A(\omega)$*

Reviews:

Optics and astrophysics: N. Wu, Springer ('97)

Spin systems: Jarrell & Gubernatis, Phys. Rep. 269 ('96)

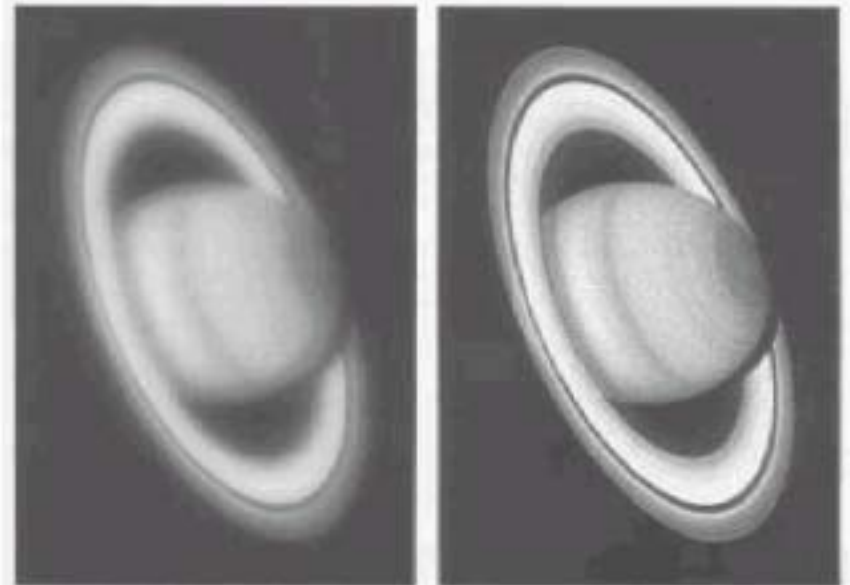
Lattice QCD: Asakawa, Nakahara and T.H., Prog. Part. Nucl. Phys. 47 ('01)

MEM Image Reconstruction

The girl's portrait

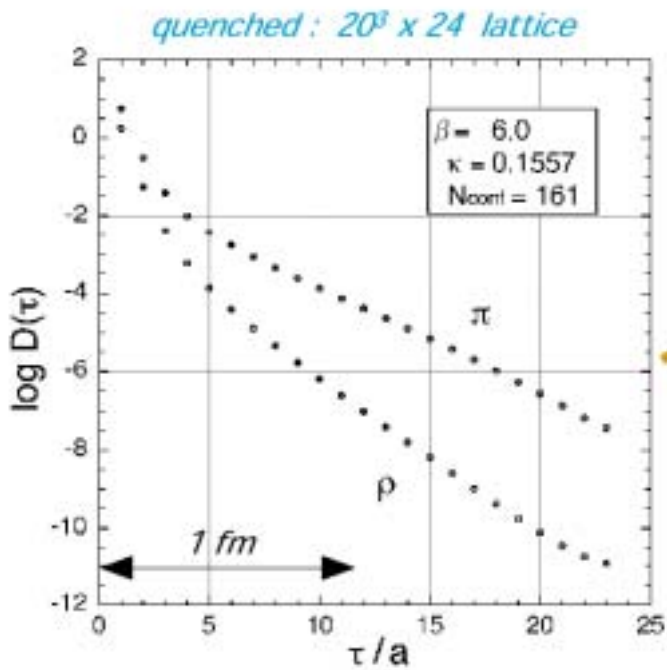


The Image of Saturn

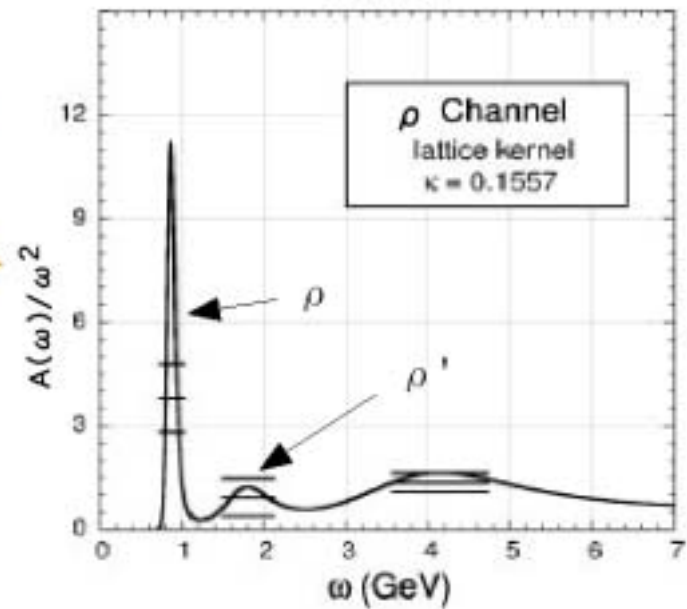
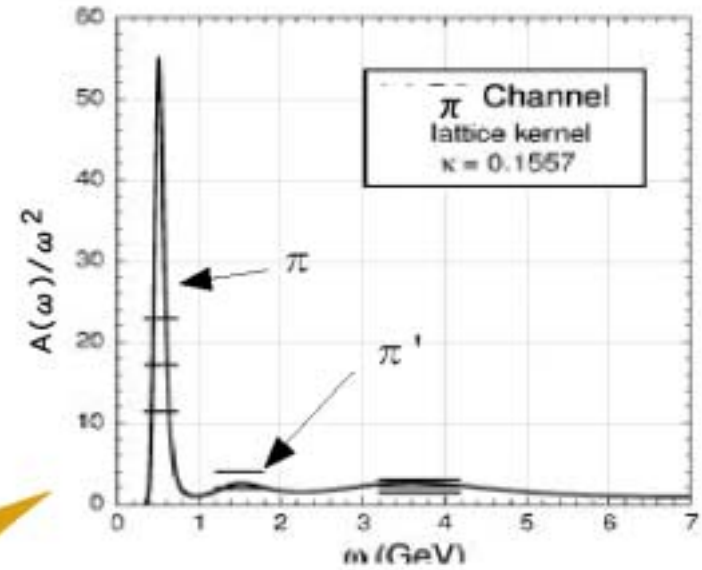


From N. Wu, "The Maximum Entropy Method" (1997)

π & ρ spectral functions
from lattice data at $T=0$



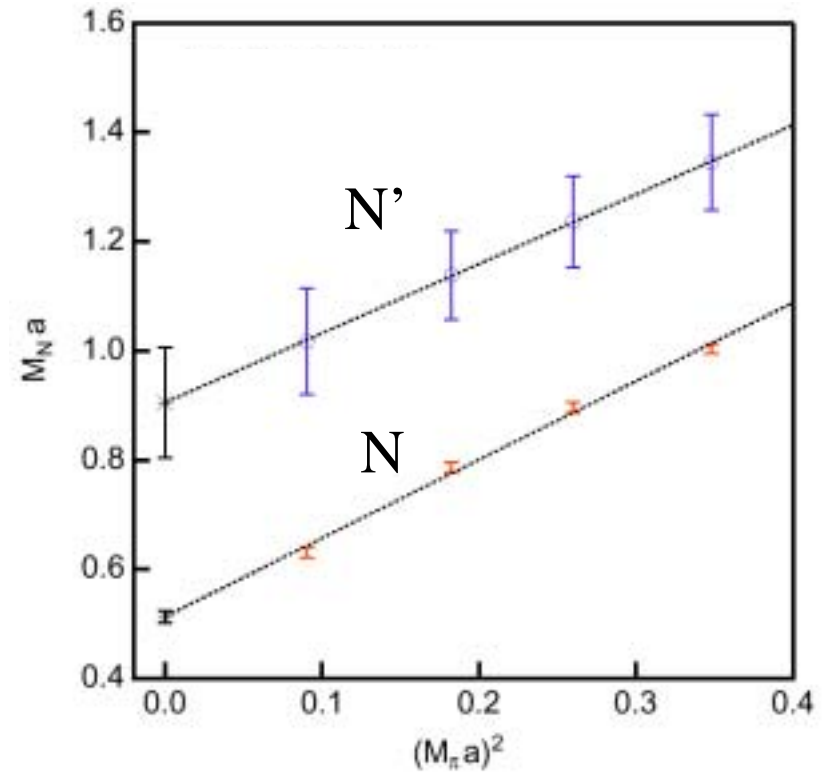
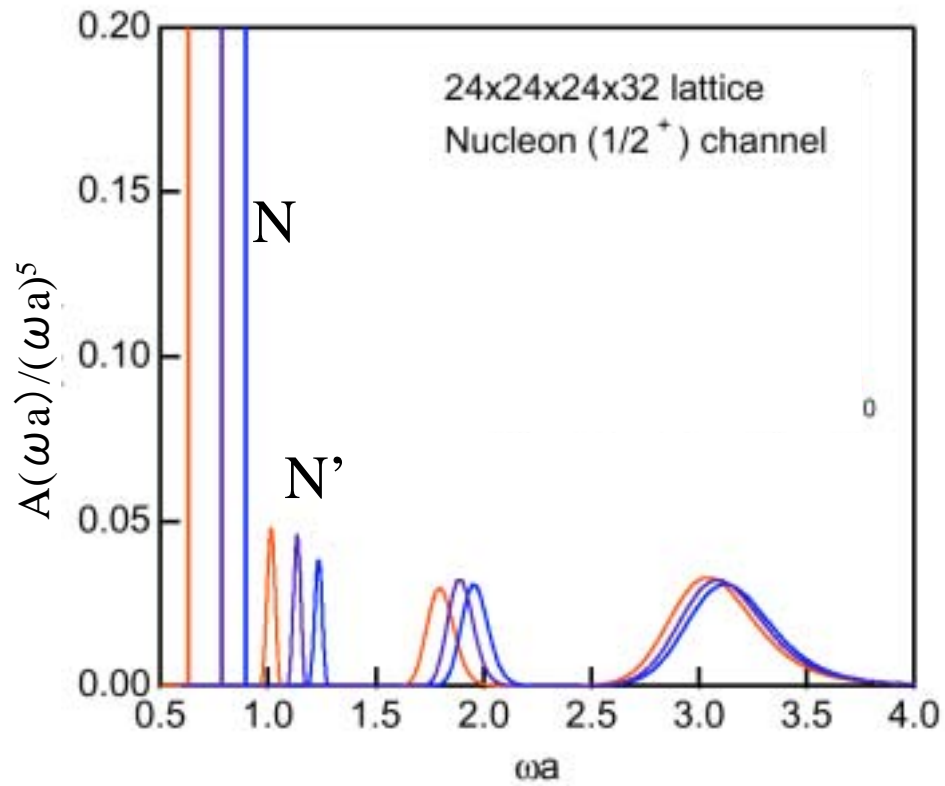
MEM



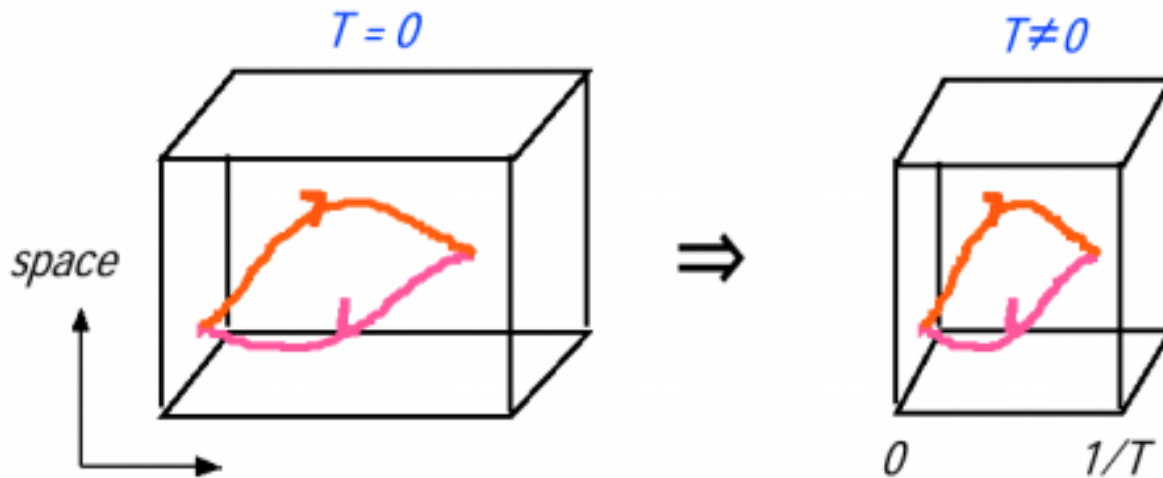
Nucleon spectral function
from lattice data at $T=0$

Sasaki, Sasaki, Asakawa + T.H.,

hep-lat/0209059 ('02)



Need of anisotropic lattice at $T \neq 0$



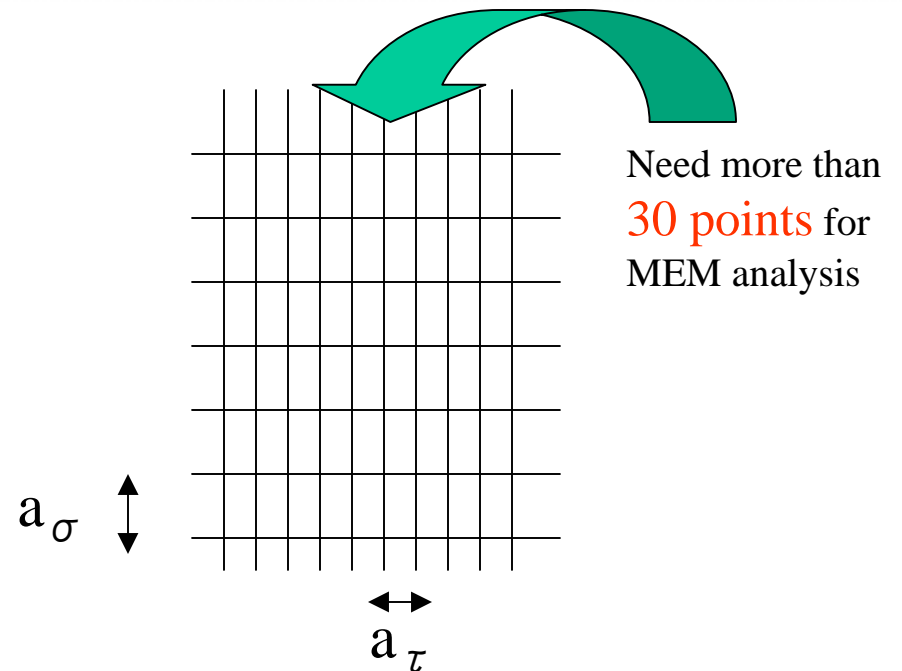
Anisotropic lattice

$$\beta = 7.0, 32^3 \times N_\tau$$

$$a_\tau = a_\sigma / 4 \approx 0.01 \text{ fm}$$

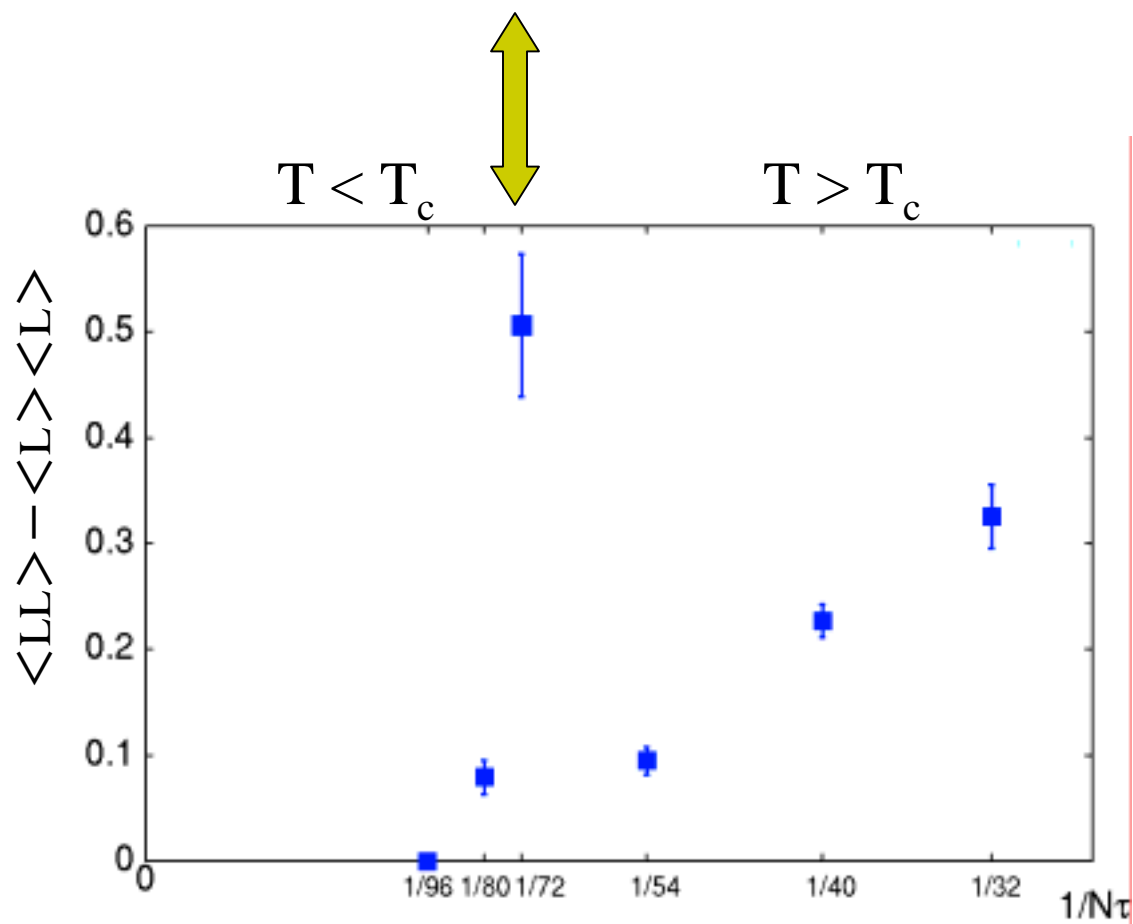
quenched approx.

Machine: CP-PACS



Temporal lattice size and T

N_τ	96	80	72	54	46	40	32
T / T_c	0.78	0.93	1.04	<u>1.4</u>	1.6	<u>1.9</u>	2.3



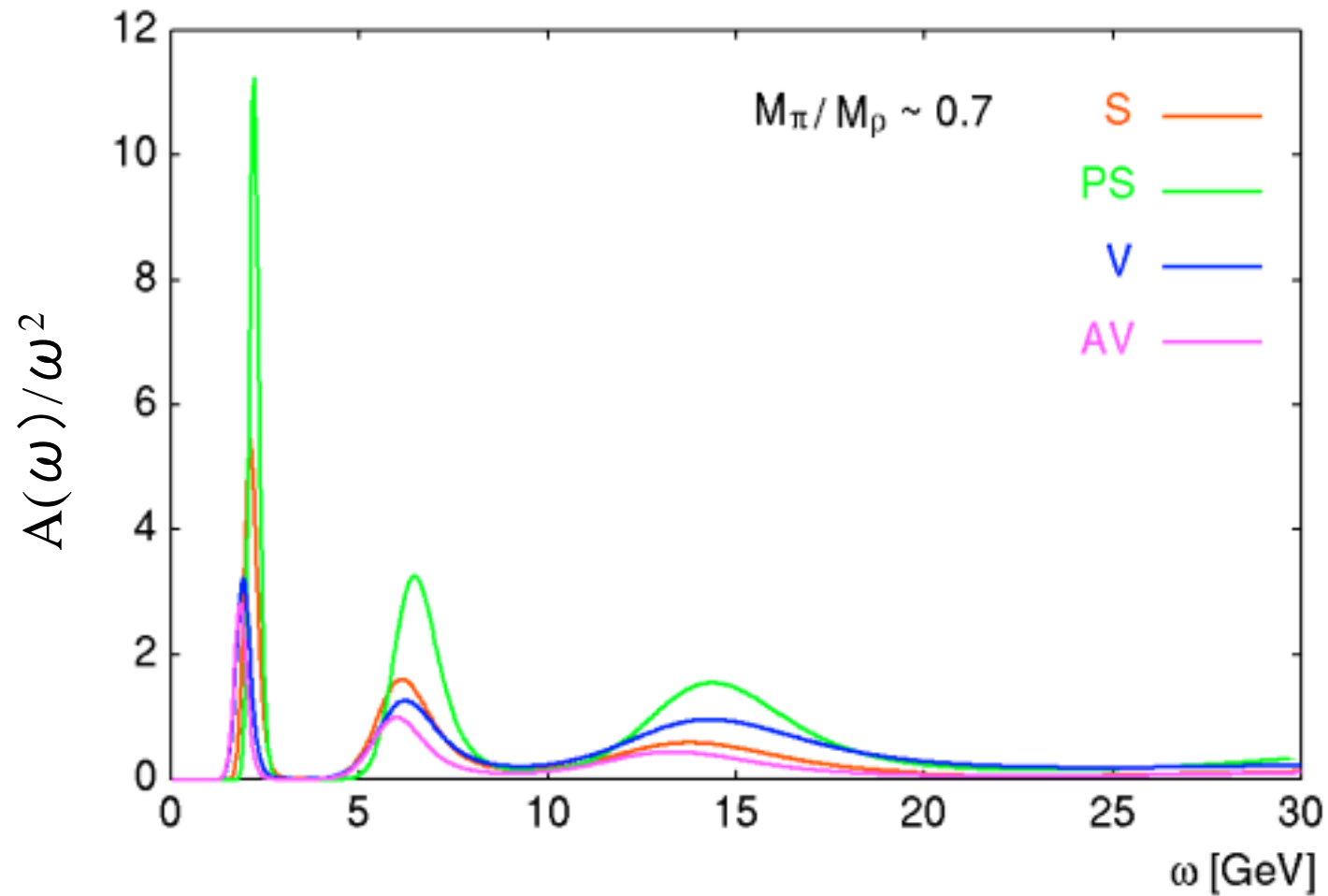
Chiral partners in the vacuum

T.H., soft-dilepton workshop at LBNL (^97)



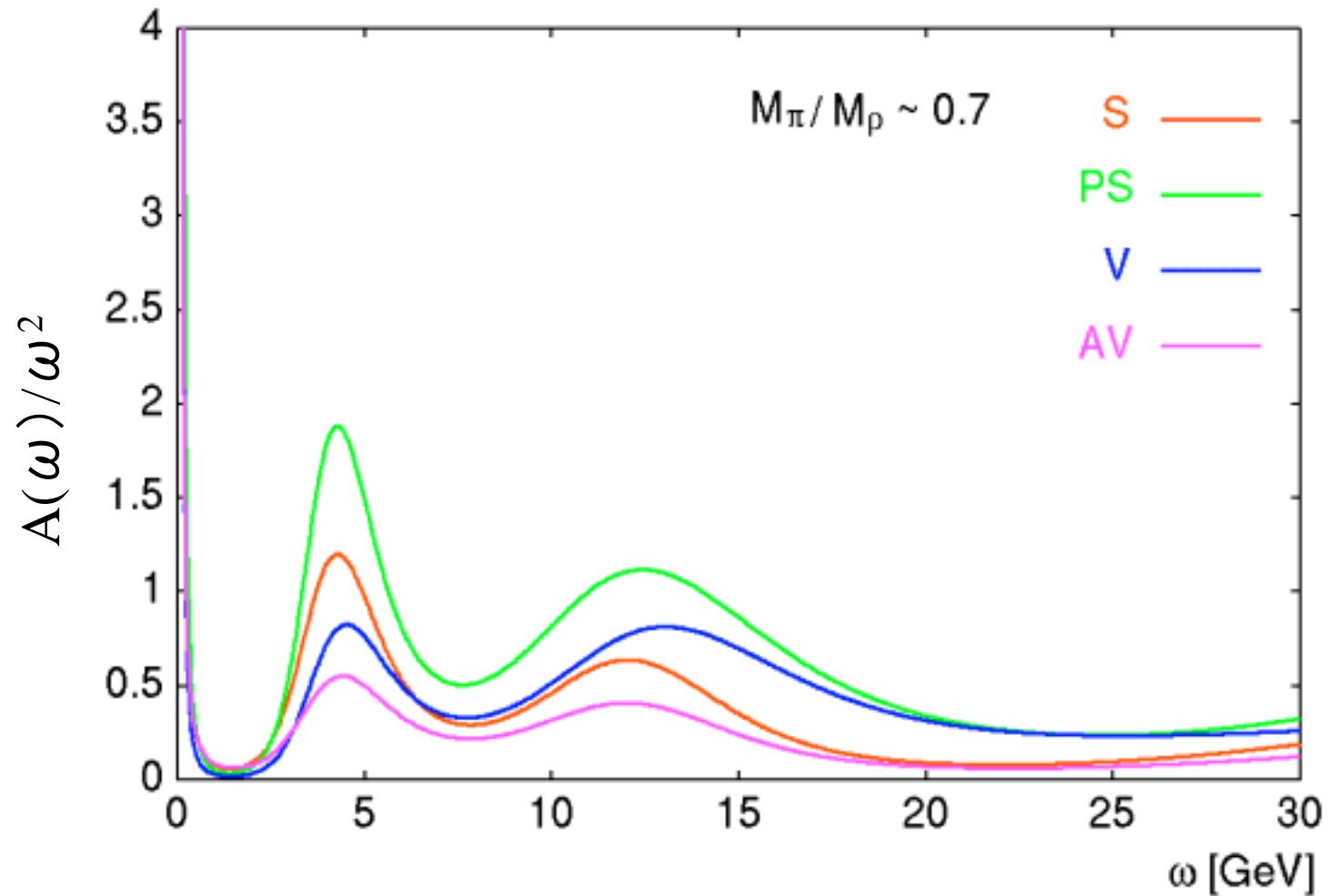
Chiral degeneracy and non-trivial modes ($T \gtrsim T_c$)

$$T = 1.4 T_c$$



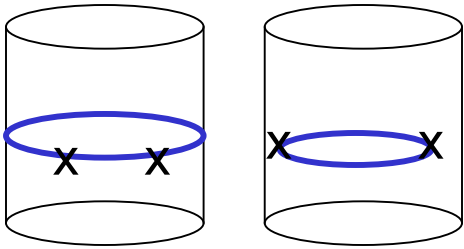
Disappearance of non-trivial modes ($T \gtrsim 2T_c$)

$$T = 1.9 T_c$$

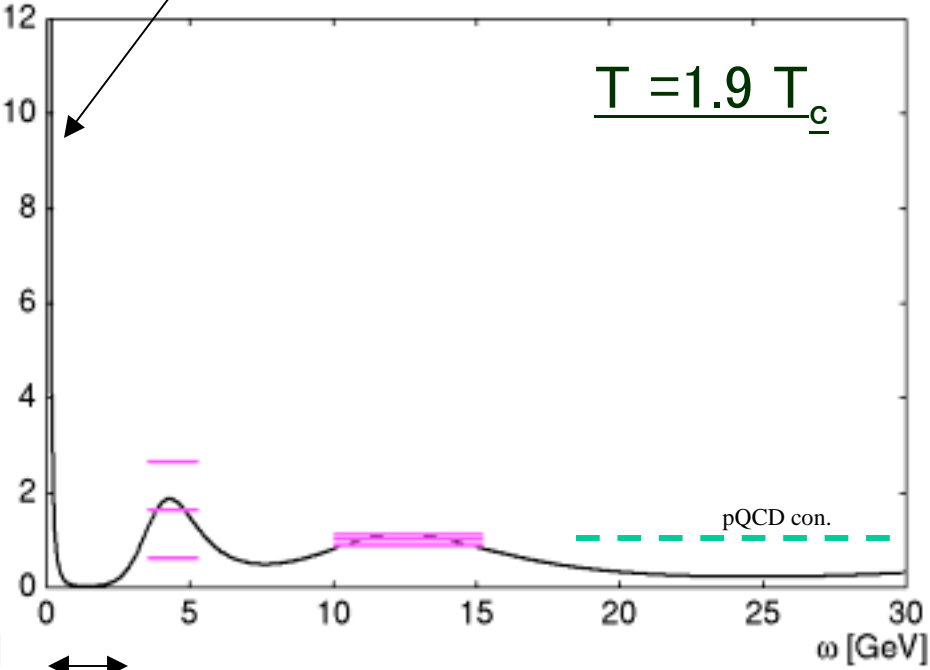
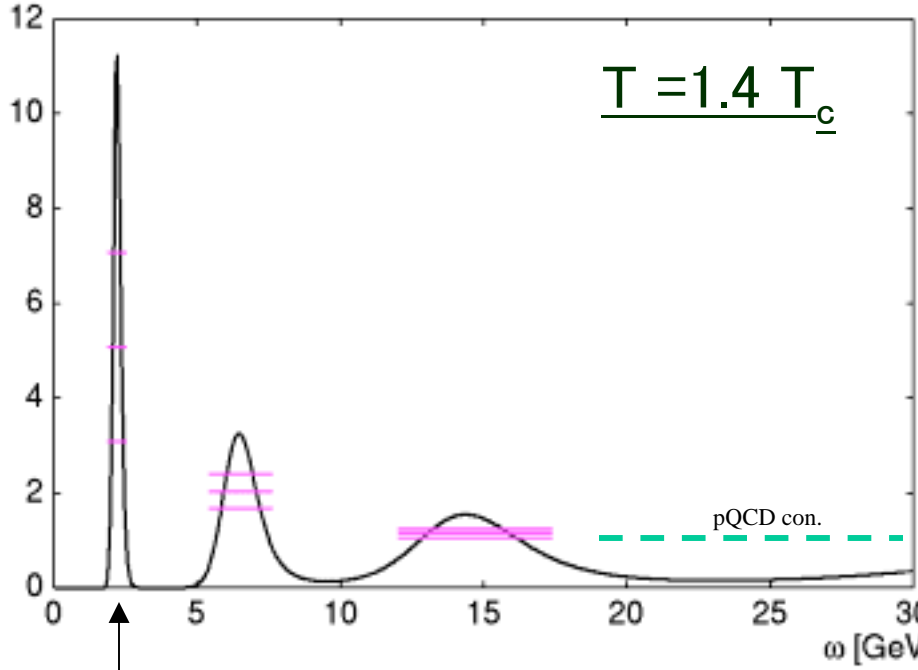


Comparison of $T = 1.4 T_c$ and $1.9 T_c$

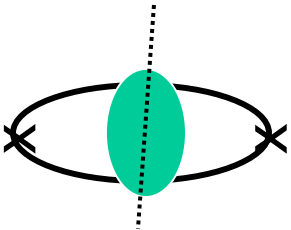
PS Channel ($M_\pi/M_\rho = 0.7$)



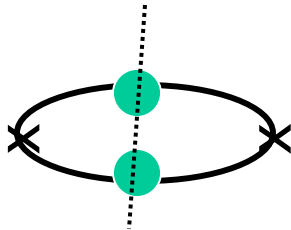
Inverse Cherenkov effect?
Koike, Lee + T.H., NPB ('93)



Para-pion mode?
Kunihiro + T.H., PRL ('85)

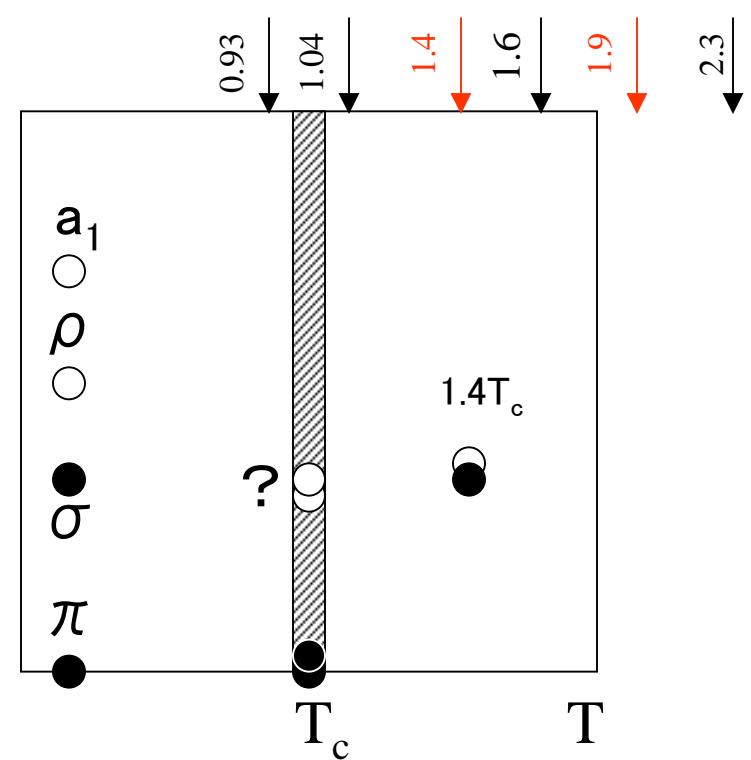
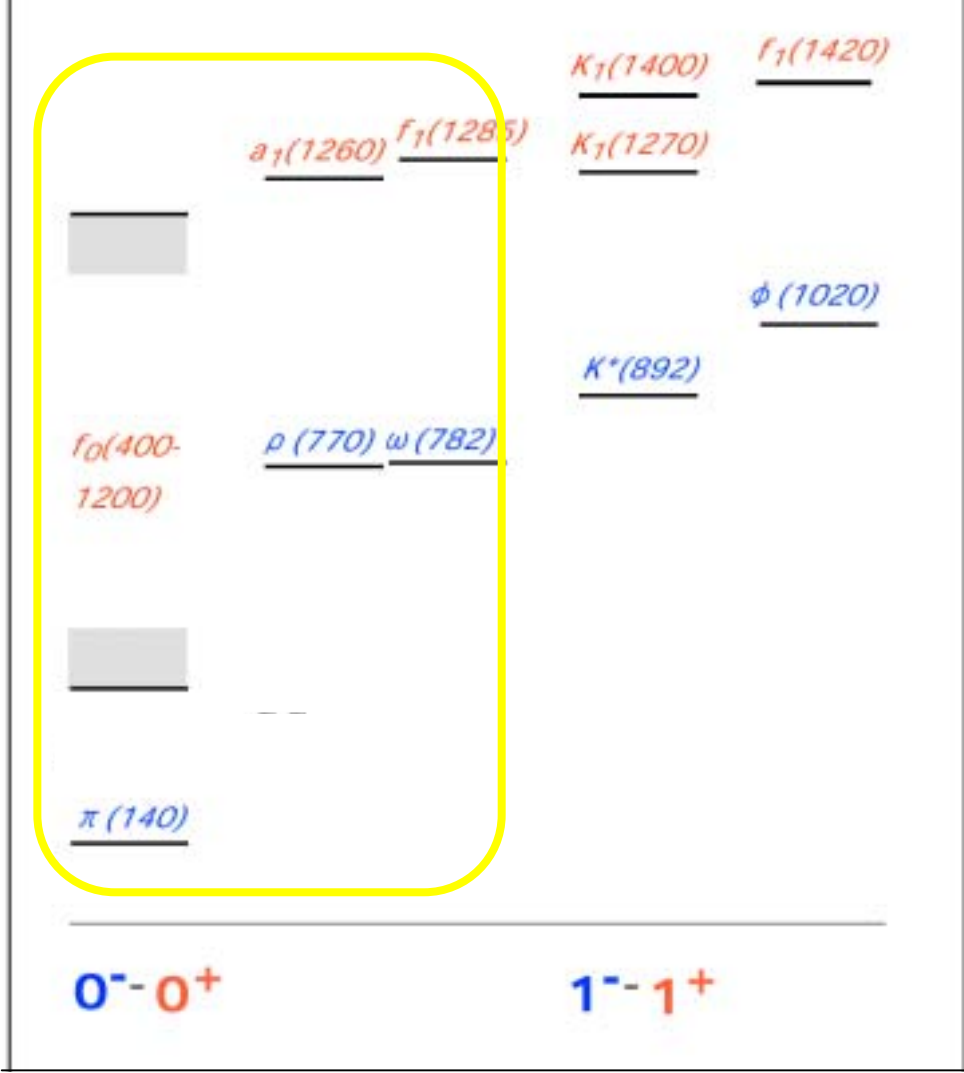


Plasmino gap?
Weldon, PRD ('82)



Hadronic modes below and above the chiral transition

T.H., soft-dilepton workshop at LBNL ('97)



Spectral function (temporal correlation):
Asakawa & T.H., hep-lat/0209059 ('02)

Summary

[1] MEM : new way of analyzing lattice QCD data

- $D(\tau) \rightarrow A(\omega)$: parameterization free, unique solution, significance test
- better lattice data \rightarrow better $A(\omega)$ Asakawa, Nakahara+T.H., (MELQCD, Tokyo)

[2] $T=0$: ground and excited states

- useful for hadron spectroscopy (π^* , ρ^* , N^* , Δ^* , glueballs etc)
Sasaki+Sasaki., (MELQCD, Tokyo)

[3] $T \neq 0$: spectral change of hadrons (thanks to CP-PACS !)

- non-trivial collective modes exist up to $2T_c$?
- evidence of the plasmino gap ?

Asakawa +T.H., (MELQCD, Tokyo)

[4] $T \neq 0$: future

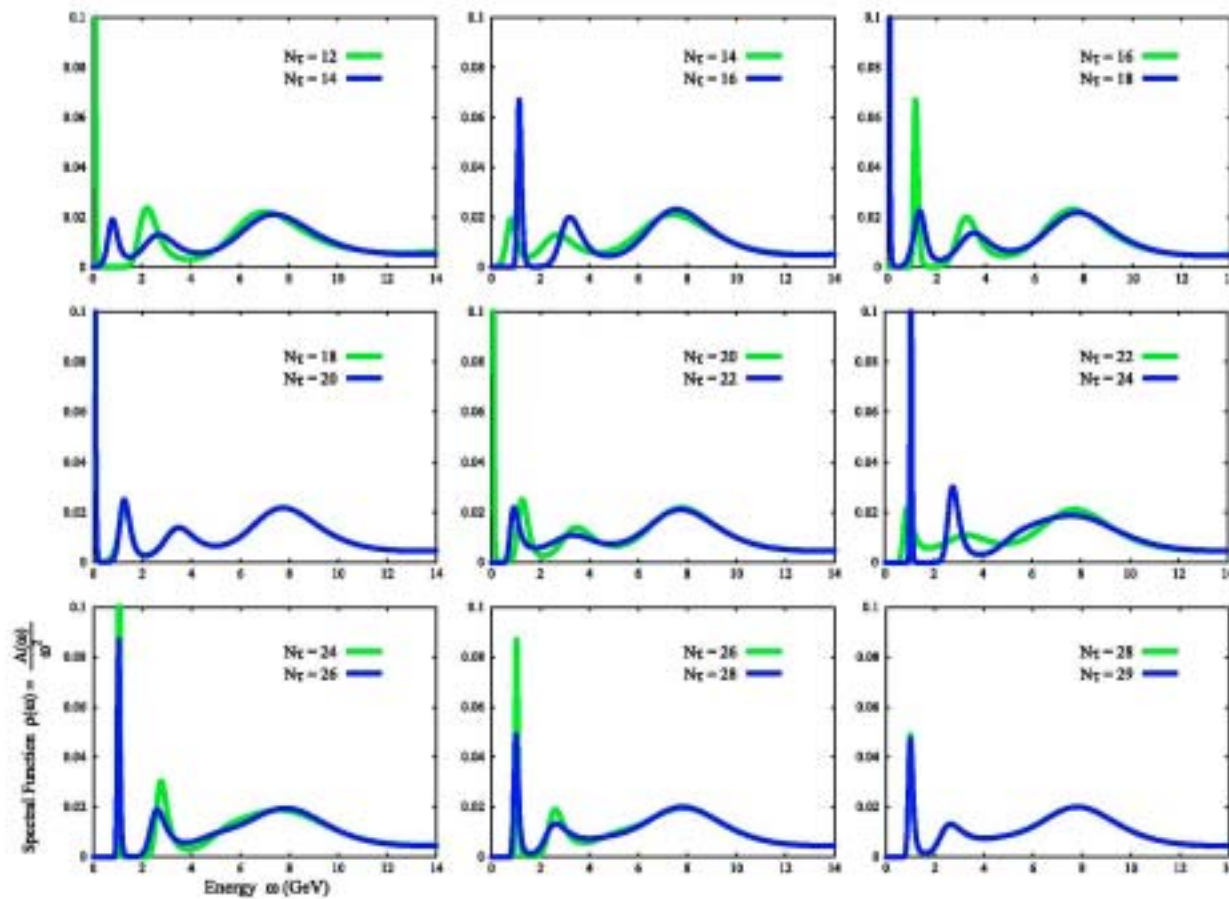
- ρ , ω , ϕ vs chiral restoration, J/ψ , ψ' vs deconfinement
- transport coefficients in hot plasma
- collective modes in $N_c=2$ dense QCD

(拡大MELQCD)

Parameters

1. Lattice size
 - $32^3 \times 32$ ($T \simeq 2.5 T_c$)
 - 40 ($T \simeq 2.0 T_c$)
 - 54 ($T \simeq 1.5 T_c$)
 - 72 ($T \simeq 1.1 T_c$)
 - 80 ($T \simeq 1.0 T_c$)
 - 96 ($T < T_c$)
2. $\beta = 7.0, \quad \xi_0 = 3.5$
3. $\xi = a_\sigma / a_\tau = 4$
 - $a_\tau = 9.75 \times 10^{-3}$ fm
 - $L_\sigma = 1.25$ fm
4. Naive Unimproved Action
5. Wilson Fermion
6. Heatbath : Overrelaxation
= 1 : 4

1000 sweeps between measurements
7. Four Quark Masses
 $m_\pi / m_\rho \simeq 0.7, 0.8, 0.9$
 $m_V = m_{J/\psi}$
8. Quenched Approximation
9. Gauge Unfixed
10. $\vec{p} = \vec{0}$ Projection
11. Machine **CP-PACS**



$40^3 \times 30$ lattice

$\beta = 6.47$

isotropic lattice



$N_\tau \simeq 30$ or larger : needed

Principles of MEM

$D(\tau) \rightarrow A(\omega)$: not unique !

What is the most probable $A(\omega)$ for given $D(\tau)$?



T. Bayes 1702-1761

Statistical Inference

Posterior prob.
likelihood func.
Prior prob.

↓
↓
↓

$$P[A|D] = \frac{1}{P[D]} P[D|A] P[A]$$

- $\frac{\delta P[A|D]}{\delta A} = 0 \Rightarrow A_{out}(\omega)$
- $\frac{\delta^2 P[A|D]}{\delta A \delta A} \Rightarrow \text{reliability of } A_{out}(\omega)$

Bayes Theorem

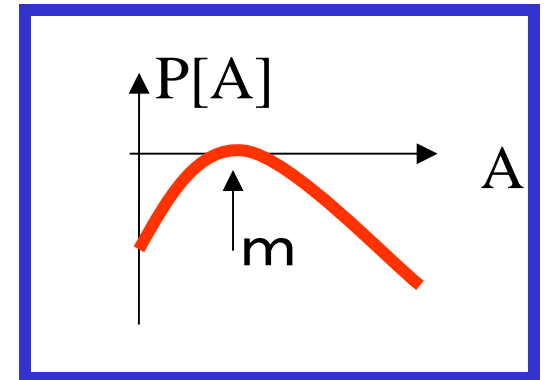
$$\begin{aligned}
 P[X|Y] &= P[X|Y] P[Y] \\
 &= P[Y|X] P[X]
 \end{aligned}$$

$P[D|A]$: Normal distribution (central limit theorem)

$P[A]$: Shanon-Jaynes information entropy

► *Explicit form of the posterior prob.*

$$P[A|D] \propto P[D|A] P[A] \propto e^{Q(D,A)}$$



• *Central limiting theorem*

$$P[D|A] \propto e^{-L(D,A)} = \exp\left(-\frac{1}{2} \sum_{i,j} [D(\tau_i) - D_A(\tau_i)] C_{ij}^{-1} [D(\tau_j) - D_A(\tau_j)]\right)$$

• *Information entropy (Shanon-Jaynes)*

$$P[A] \propto e^{\alpha S(A)} = \exp \alpha \int_0^{\infty} \left[A(\omega) - m(\omega) - A(\omega) \ln \left(\frac{A(\omega)}{m(\omega)} \right) \right] d\omega$$

- ◆ $Q(D,A) = \alpha S - L$: "- free energy" ⇐ *to be maximized with respect to $A(\omega)$*
- ◆ α : a fictitious "temperature" ⇐ *to be integrated with a weight $P[\alpha | D]$*
- ◆ $m(\omega)$: prior estimate of $A(\omega)$ ⇐ *to be updated by error analysis*

► Procedure of modern MEM

Step 1 Maximizing Q

$$\frac{\delta Q}{\delta A(\omega)} = 0 \longrightarrow A_{\alpha}(\omega)$$

- Solution is unique (Asakawa, T.H., Nakahara (00))
- Rapid convergence by SVD (Bryan, EBJ (90))

Step 2 averaging over α

$$A_{out}(\omega) = \int A_{\alpha}(\omega) P[\alpha | D] d\alpha$$

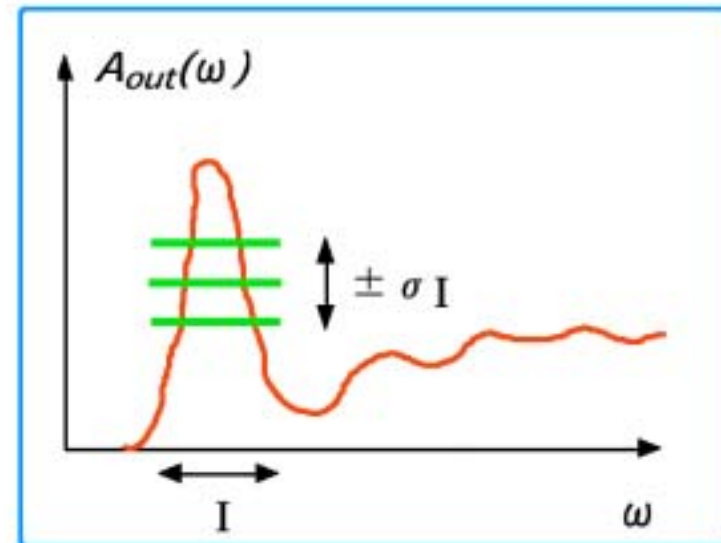
with

$$P[\alpha | D] \sim \int [dA] P[D|A\alpha] P[A|\alpha] P[\alpha]$$

$$\ln P[\alpha | D] = const. + \frac{1}{2} \sum_k \ln \frac{\alpha}{\alpha + \lambda_k} + Q(D, A_{\alpha})$$

Step 3 error analysis

$$\sigma_I^2 = - \left\langle \left(\frac{\delta^2 Q}{\delta A(\omega) \delta A(\omega')} \right)_{A=A_{\alpha}}^{-1} \right\rangle_I$$

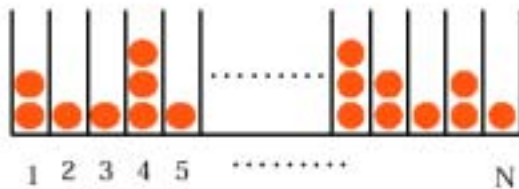


Shannon-Jaynes Entropy

- Combinatorial construction (Monkey argument)

*Friedan (72); Gull & Daniell (79);
Jaynes (86); Skilling (88)*

Product of Poisson distribution → SJ entropy



$$P_{\lambda}(n) = \prod_{i=1}^M \frac{\lambda_i^{n_i} e^{-\lambda_i}}{n_i!}$$

$$\rightarrow \exp \left[\alpha \sum_{i=1}^M \left(A_i - m_i - A_i \ln \left(\frac{A_i}{m_i} \right) \right) \right]$$

$$A_i = n_i / \alpha$$

$$m_i = \lambda_i / \alpha$$

$$n! \sim \exp [n \log n - n]$$

- Axiomatic construction

Axiom I: Locality

Axiom II: Coordinate invariance

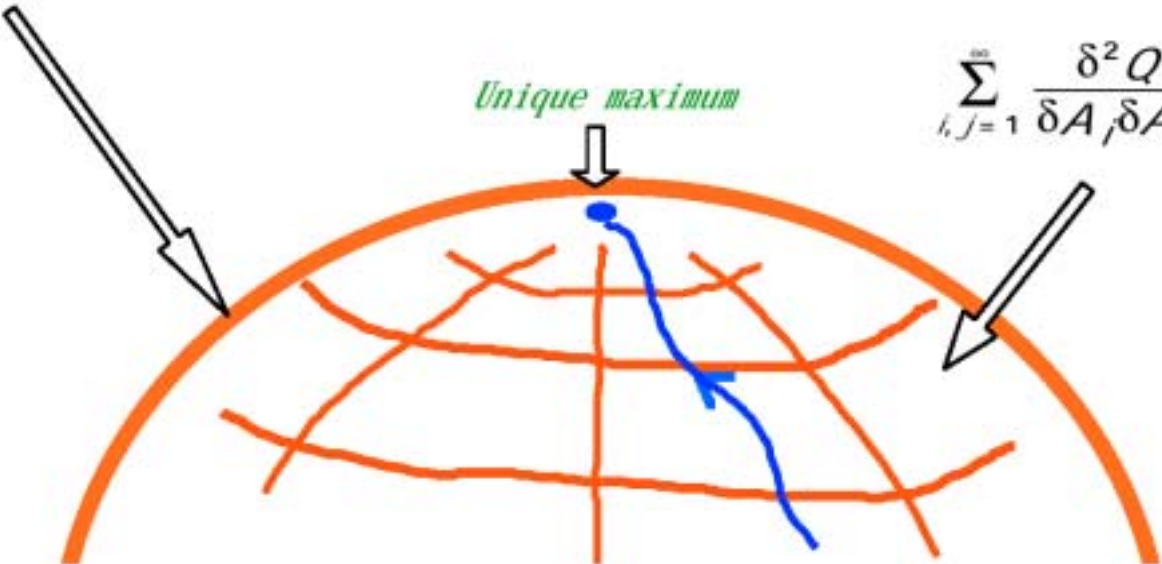
Axiom III: System independence

Axiom IV: Scaling

*Shanon (1948); Jaynes (1957), ...,
Skilling (88);
Asakawa, T.H., & Nakahara, (01)*

Properties of hypersurface $Q(A)$

$Q(A_i) : 0(1000)$ -dim. surface



$$\sum_{i,j=1}^{\infty} \frac{\delta^2 Q}{\delta A_i \delta A_j} z_i z_j > 0$$

$$\frac{\delta Q}{\delta A_i} = 0$$

$N_{data} = 0(10)$ -dim. space

Brvan, 1990

► Kernel $K(\tau, \omega)$

Free propagator with mass ω

$$D(\tau > 0) = \int_0^{\infty} K(\tau, \omega) A(\omega) d\omega$$

• "continuum" kernel

$$K_{cont}(\tau, \omega) = e^{-\omega \tau} = 2\omega \int_{-\infty}^{\infty} \frac{d\nu}{2\pi} \frac{e^{i\nu\tau}}{\omega^2 + \nu^2} d\nu$$

• "lattice" kernel

$$K_{lat}(\tau, \omega) \equiv 2\omega \int_{-\pi/a}^{\pi/a} \frac{d\nu}{2\pi} \frac{e^{i\nu\tau}}{\omega^2 + \left(\frac{2}{a} \sin \frac{\nu a}{2}\right)^2} d\nu$$

$O(a)$ difference